## Conclusion

From the data presented in this work we conclude that quasiergodicity is reached relatively quickly and that the used equilibration times are adequate. For bulk water and the solvated calcium ion, static properties seem to be plausibly represented by the present potential, as are also qualitatively dynamic properties. The failure to quantitatively reproduce dynamic, experimental data is mainly due to the absence of dynamic data in constructing the intermolecular potential. A water model that allows for intramolecular motion while correctly reproducing dynamic properties is desirable but lacking. Also for the larger EDTA molecule, structural properties are qualitatively reproduced but are more sensitive to the imperfections of the potential. The binding of the calcium ion to EDTA is adequately treated in most
respects. The CaEDTA ${ }^{2-}$ overall mobility is too large, but this may be corrected for by use of the known properties of the water model and experimental data. While shortcomings of the intramolecular EDTA potential render the simulation dynamics locally incorrect, the mechanism of this seems well understood and can be taken into account in the interpretation of simulations. We think the present simulations contribute substantially to the understanding of calcium binding, which; in view of its biological importance, certainly merits further investigation.

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Registry No. $\mathrm{Ca}, 7440-70-2 ; \mathrm{H}_{2} \mathrm{O}, 7732-18-5$; EDTA, 60-00-4.

# Calculated Structures and Fluoride Affinities for Fluorides 

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#### Abstract

It is shown that SCF-MO calculations provide good estimates of the energies of the processes $\mathrm{MF}_{n} \rightarrow \mathrm{M}^{n+}+n \mathrm{~F}^{-}$ where $\mathrm{M}^{n+}$ is an ion of a first- or second-row element in a closed-shell or $\mathrm{s}^{2}$ configuration. The fluoride ion affinities are then calculated for a number of molecules and ions. Where comparison with experiment is possible, the agreement is generally good when allowance is made for experimental uncertainties. In favorable cases, accurate heats of formation may be calculated from fluoride affinities.


Fluoride ion affinities of molecules and ions are of considerable theoretical and experimental interest. Experimental values come mainly from ion cyclotron resonance (ICR) experiments ${ }^{1-3}$ or from Born-Haber ( BH ) thermodynamic cycles. ${ }^{4,5}$ Most of the derived results are indirect, and it is often quite difficult to assess uncertainties reliably. It is important therefore to have an independent set of fluoride affinities to compare with experimental values and to provide estimates of values as yet not determined.

In this work it is shown that ab initio SCF-MO calculations can provide reliable fluoride affinities, at least in certain instances. The important point about such calculations is that they should be made only for processes in which there is no change in multiplicity and ideally for singlet state species. ${ }^{6}$ Previous calculations of the fluoride affinity of $\mathrm{HF}^{7,8}$ and of $\mathrm{H}_{2} \mathrm{O}^{9}$ have shown that in these cases at least the change of correlation energy is only a small fraction of the energy change and less than experimental error, which is typically $>5 \mathrm{~kJ} \mathrm{~mol}^{-1}$.

## Methods

Equilibrium geometries and SCF energies have been calculated for a number of first- and second-row fluoride molecules and ions in singlet states. The $6-31 \mathrm{G}^{*}$ basis set ${ }^{10}$ was chosen as a reasonable

[^0]Table I. Energy ( $\mathrm{kJ} \mathrm{mol}^{-1}$ ) for $\mathrm{MF}_{n} \rightarrow \mathrm{M}^{n+}+n F^{-}$

| $\mathrm{MF}_{n}$ | expt $^{a}$ | theor | ratio $^{b}$ |
| :--- | ---: | ---: | :--- |
| HF | 1573 | 1591 | 0.989 |
| LiF | 772 | 764 | 1.011 |
| $\mathrm{BeF}_{2}$ | 3285 | 3273 | 1.004 |
| $\mathrm{BF}_{3}$ | 1241 | 1196 | 1.038 |
| $\mathrm{BF}_{3}$ | 7859 | 7800 | 1.008 |
| $\mathrm{CF}_{2}{ }^{c}$ | 3848 | 3762 | 1.023 |
| $\mathrm{CF}_{4}$ | 14949 | 14775 | 1.012 |
| $\mathrm{NF}_{3}$ | 8700 | 8510 | 1.022 |
| $\mathrm{NaF}^{\mathrm{MgF}}$ |  | 650 | 635 |
| $\mathrm{AlF}_{2}$ | 2564 | 2533 | 1.024 |
| $\mathrm{AF}_{3}$ | 913 | 906 | 1.012 |
| $\mathrm{SiF}_{2}$ | 5931 | 5869 | 1.008 |
| $\mathrm{SiF}_{4}$ | 11036 | 2866 | 1.011 |
| $\mathrm{PF}_{3}$ | 6369 | 10863 | 1.016 |
| $\mathrm{PF}_{5}$ | 17789 | 6191 | 1.029 |
| $\mathrm{SF}_{4}$ | 11272 | 17421 | 1.021 |
| $\mathrm{SF}_{6}$ | 26701 | 10917 | 1.033 |

${ }^{a}$ Data from ref 13 except electron affinity of $F$ from ref 14 and zero point energies estimated from frequencies given in ref 15 . ${ }^{b} \mathbf{E x p t} /$ theor. ${ }^{c}$ Heat of formation from ref 15 .
compromise between the desire, on the one hand, to have a sufficiently flexible basis set to predict geometries reliably and to avoid basis set superposition errors and, on the other hand, to allow reasonably heavy molecules to be studied without untoward cost. It is well-known ${ }^{11}$ that for negative ions ( $\mathrm{F}^{-}$, etc.) additional diffuse valence orbitals are necessary, so the F basis was supplemented

[^1]Table II. Calculated Geometries and Energies ${ }^{a}$

| species | symmetry | energy | dimensions | $E_{\mathrm{F}}$ | species | symmetry | energy | dimensions | $E_{\text {F }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| F | $K_{h}$ | -99.4178 |  |  | $\mathrm{AlF}_{3}$ | $D_{3 h}$ | -540.4677 | 1.630 (1.63) | 510 |
| HF | $C_{\infty 0}$ | -100.0237 | 0.902 (0.917) | 177 | $\mathrm{AlF}_{4}{ }^{-}$ | $T_{d}$ | -640.0778 | 1.690 | -243 |
| $\mathrm{HF}_{2}{ }^{-}$ | $D_{\text {ch }}$ | -199.5098 | 1.125 |  | $\mathrm{AlF}_{5}{ }^{2-}$ | $D_{3 h}$ | -739.4043 | $1.814,{ }^{\text {b }} 1.768$ | -680 |
| $\mathrm{Li}^{+}$ | $K_{h}$ | -7.2355 |  | 774 | $\mathrm{AlF}_{6}{ }^{\text {- }}$ | $O_{h}$ | -838.5658 | 1.883 |  |
| LiF | $C_{\infty \nu}$ | -106.9443 | 1.589 (1.564) | 309 | $\mathrm{Al}^{+}$ | $K_{h}$ | -241.6528 |  | 913 |
| $\mathrm{LiF}_{2}{ }^{-}$ | $D_{\text {mh }}$ | -206.4786 | 1.695 |  | AIF | $C_{\infty}$ | -341.4156 | 1.670 | 306 |
| $\mathrm{Be}^{2+}$ | $K_{h}$ | -13.6098 |  | 2009 | $\mathrm{AlF}_{2}{ }^{\text {- }}$ | $C_{2 v}$ | -440.9491 | $1.735,99.3$ |  |
| $\mathrm{BeF}^{+}$ | $C_{\infty}$ | -113.7903 | 1.321 | 1276 | $\mathrm{Si}^{4+}$ | $K_{h}$ | -285.1563 |  | 4302 |
| $\mathrm{BeF}_{2}$ | $D_{\text {oh }}$ | -213.6921 | 1.378 | 391 | $\mathrm{SiF}^{3+}$ | $C_{\text {cu }}$ | -386.1869 | 1.524 | 3274 |
| $\mathrm{BeF}_{3}{ }^{-}$ | $D_{3 h}$ | -313.2586 | 1.474 | -256 | $\mathrm{SiF}_{2}{ }^{+}$ | $D_{\text {en }}$ | -486.8321 | 1.485 | 2167 |
| $\mathrm{BeF}_{4}{ }^{2-}$ | $T_{d}$ | -412.5785 | 1.597 |  | $\mathrm{SiF}_{3}{ }^{+}$ | $D_{3 h}$ | -587.0623 | 1.515 | 1294 |
| Be | $K_{h}$ | -14.5669 |  | 276 | $\mathrm{SiF}_{4}$ | $T_{d}$ | -686.9652 | 1.561 (1.554) | 323 |
| $\mathrm{BeF}^{-}$ | $C_{\infty}$ | -114.0894 | 1.454 |  | $\mathrm{SiF}_{5}{ }^{-}$ | $D_{3 h}$ | -786.5042 | $1.663,{ }^{\text {b }}, 1.623$ | -187 |
| $\mathrm{B}^{3+}$ | $K_{h}$ | -21.9839 |  | 3930 | $\mathrm{SiF}_{6}{ }^{2-}$ | $O_{h}$ | -885.8519 | 1.710 (1.708) |  |
| $\mathrm{BF}^{2+}$ | $C_{\text {mo }}$ | -122.8868 | 1.176 | 2671 | $\mathrm{Si}^{2+}$ | $K_{h}$ | -287.9704 |  | 1896 |
| $\mathrm{BF}_{2}{ }^{+}$ | $D_{\text {en }}$ | -223.3138 | 1.218 | 1262 | $\mathrm{SiF}^{+}$ | $C_{\infty}$ | -388.0914 | 1.539 | 1048 |
| $\mathrm{BF}_{3}$ | $D_{3 h}$ | -323.2083 | 1.303 (1.313) | 358 | $\mathrm{SiF}_{2}$ | $C_{2 v}$ | -487.8976 | 1.604, 99.2 | 366 |
| $\mathrm{BF}_{4}{ }^{-}$ | $T_{d}$ | -422.7612 | 1.397 |  |  |  |  | (1.590, 100.8) |  |
| $\mathrm{B}^{+}$ | $K_{h}$ | -24.2341 |  | 1241 | $\mathrm{SiF}_{3}{ }^{-}$ | $\mathrm{C}_{30}$ | -587.4506 | 1.673, 96.1 |  |
| BF | $C_{\text {me }}$ | -124.1074 | 1.266 (1.265) | 187 | $\mathrm{P}^{\text {S }}$ | $K_{h}$ | -334.3411 |  |  |
| $\mathrm{BF}_{2}{ }^{\text {+ }}$ | $C_{20}$ | -223.5936 | 1.422, 100.6 |  | $\mathrm{PF}_{2}{ }^{3+}$ | $D_{\text {oh }}$ | -537.3863 | 1.437 | 3310 |
| $\mathrm{CF}^{4+}$ | $K_{h}$ | -32.3577 -1347377 |  | 6913 | $\mathrm{PF}_{3}{ }^{2+}$ | $D_{3 h}$ | -638.0389 | 1.443 | 2192 |
| $\mathrm{CF}^{3+}{ }^{\text {+ }}$ | $C_{\text {cov }}$ | -134.3773 | 1.112 | 4559 | $\mathrm{PF}_{4}{ }^{+}$ | $T_{d}$ | -738.2785 | 1.480 | 1561 |
| $\mathrm{CF}_{2}{ }^{2+}$ | $D_{\infty}$ | -235.5110 | 1.134 | 2344 | $\mathrm{PF}_{5}$ | $D_{3 h}$ | -838.0654 | $1.572,61.536$ | 406 |
| $\mathrm{CF}_{3}{ }^{+}$ | $D_{3 h}$ | -335.8109 | 1.218 | 1136 |  |  |  | $(1.577,1.534)$ |  |
| $\mathrm{CF}_{4}$ | $T_{d}$ | -435.6565 | 1.304 (1.320) |  | $\mathrm{PF}_{6}{ }^{-}$ | $O_{h}$ | -937.6346 | 1.610 |  |
| $\mathrm{C}^{2+}$ | $K_{h}$ | -36.3992 |  | 2767 | $\mathrm{P}^{3+}$ | $K_{h}$ | -338.5328 |  | 3220 |
| $\mathrm{CF}^{+}$ | $C_{\text {cu }}$ | -136.8472 | 1.146 | 1081 | $\mathrm{PF}^{2+}$ | $C_{\infty}$ | -439.1424 | 1.442 | 2022 |
| $\mathrm{CF}_{2}$ | $C_{2 v}$ | -236.6677 | $\begin{gathered} 1.288,104.2 \\ (1.300,104.9) \end{gathered}$ | 169 | $\begin{aligned} & \mathrm{PF}_{2}{ }^{+} \\ & \mathrm{PF}_{3} \end{aligned}$ | $\begin{aligned} & C_{2 v} \\ & C_{3 v} \end{aligned}$ | $\begin{aligned} & -539.3086 \\ & -639.1442 \end{aligned}$ | $\begin{aligned} & 1.504,100.9 \\ & 1.570,97.1 \end{aligned}$ | 1129 189 |
| $\mathrm{CF}_{3}{ }^{-}$ | $\mathrm{C}_{30}$ | -336.1484 | 1.420, 99.1 |  |  |  |  | (1.570, 97.8) |  |
| $\mathrm{N}^{3+}$ | $K_{h}$ | -51.0574 |  | 5247 | $\mathrm{PF}_{4}{ }^{-}$ | $C_{2 v}$ | -738.6321 | $1.751,167.8$ |  |
| $\mathrm{NF}^{2+}$ | $C_{\text {© }}$ | -152.4309 | 1.068 | 2367 |  |  |  | 1.611, 99.8 |  |
| $\mathrm{NF}_{2}{ }^{+}$ | $C_{2 b}$ | -252.7310 | $1.212,108.4$ | 1082 | $\mathrm{S}^{\mathbf{6 +}} \mathrm{SF}^{3+}$ | $K_{h}$ | -387.5316 |  |  |
| $\mathrm{NF}_{3}$ | $C_{30}$ | -352.5523 | $\begin{gathered} 1.330,102.7 \\ (1.365,102.4) \end{gathered}$ |  | $\begin{aligned} & \mathrm{SF}_{3}{ }^{3+} \\ & \mathrm{SF}_{4}{ }^{2+} \end{aligned}$ | $D_{3 h}$ $T_{d}$ | -693.4955 -794.1071 | 1.415 1.439 | 3200 1590 |
| $\mathrm{Na}^{+}$ | $K_{h}$ | -161.6595 |  | 650 | $\mathrm{SF}_{5}{ }^{+}$ | $D_{3 h}$ | -894.1179 | $1.524,{ }^{6} 1.492$ | 1252 |
| NaF | $C_{\text {mo }}$ | -261.3192 | 1.933 (1.926) | 273 | $\mathrm{SF}_{6}$ | $O_{h}$ | -994.0028 | 1.557 (1.564) |  |
| $\mathrm{NaF}_{2}{ }^{-}$ | $D_{\text {mh }}$ | -360.8393 | $2.036$ |  | $\mathrm{S}^{4+}$ | $K_{h}$ | -393.3312 |  |  |
| $\mathrm{Mg}^{2+}$ | $K_{h}$ | -198.8119 |  | 1526 | $\mathrm{SF}_{3}{ }^{+}$ | $C_{30}$ | -695.4164 | 1.501, 99.0 | 885 |
| $\mathrm{MgF}^{+}$ | $C_{\text {cov }}$ | -298.8041 | 1.693 | 1037 | $\mathrm{SF}_{4}$ | $C_{2 v}$ | -795.1603 | 1.643, 170.1 | 211 |
| $\mathrm{MgF}_{2}$ | $D_{\infty h}$ | -398.6124 | 1.742 | 404 |  |  |  | $(1.646,173.1)$ |  |
| $\mathrm{MgF}_{3}{ }^{-}$ | $D_{3 h}$ | -498.1821 | 1.809 | -128 |  |  |  | 1.544, 102.5 |  |
| $\mathrm{MgF}_{4}{ }^{\text {- }}$ | $T_{d}$ | -597.5521 | 1.905 |  |  |  |  | $(1.545,101.6)$ |  |
| $\mathrm{Mg}^{\mathrm{MgF}}$ | $K_{h}$ | -199.3243 |  | 956 | $\mathrm{SF}_{5}{ }^{-}$ | $C_{40}$ | -894.6560 | $1.588,{ }^{\text {c }} 84.4{ }^{\text {d }}$ |  |
| $\mathrm{MgF}^{-}$ | $C_{\text {cov }}$ | -299.1022 | 1.838 |  |  |  |  | 1.711, 89.5 |  |
| $\mathrm{Al}^{3+}{ }^{\text {+ }}$ | $K_{h}$ | -239.9789 |  | 2708 |  |  |  |  |  |
| $\mathrm{AlF}^{2+}$ | $C_{\text {cov }}$ | -340.4171 | 1.564 | 2029 |  |  |  |  |  |
| $\mathrm{AlF}_{2}{ }^{+}$ | $D_{\text {oh }}$ | -440.5995 | 1.585 | 1196 |  |  |  |  |  |

${ }^{a}$ Experimental values ${ }^{e}$ in parentheses with distances in angstroms and angles between equivalent bonds in degrees and energies in hartrees for fluorides, and fluoride affinities, $E_{\mathrm{F}}$ (in $\mathrm{kJ} \mathrm{mol}^{-1}$ ). ${ }^{b}$ Axial. ${ }^{〔}$ Apical. ${ }^{d}$ Angle between apical and basal bonds. ${ }^{e}$ References 15 and 17 .
with an additional p gaussian orbital with exponent ( 0.090 au ) optimized for $\mathrm{F}^{-}$. A similarly augmented F basis was found to give a good account of the energetics of formation of oligomers of LiF and $\mathrm{NaF} .{ }^{12}$

Thermochemical data were taken from the revision of NBS circular $500^{13}$ except that I have taken the heat of formation of $\mathrm{F}^{-}$to be $\Delta H^{\circ}=-251.1 \mathrm{~kJ} \mathrm{~mol}{ }^{-1}$ calculated from the revised electron affinity of fluorine ( 3.399 eV ) reported by Milstein and Berry. ${ }^{14}$ Zero point energies were calculated from the frequencies listed in the JANAF thermochemical tables. ${ }^{\text {is }}$

## Results

The results are listed in Tables I and II. In Table I I compare the calculated energies for $\mathrm{MF}_{n} \rightarrow \mathrm{M}^{n+}$ with thermochemical data (corrected for zero point energy). The agreement can be seen to be very good, with the calculated values generally smaller by

[^2]about $2 \%$ and no systematic variation across the periodic table. Indeed, for smaller values of $n$ the absolute error is small, of the order often found for experimental uncertainties. As the heats of formation of atomic ions are very well-known, it can be seen that the data could be used to calculate heats of formation of molecules in much the same way as proposed by Hurley. ${ }^{6}$

The origin of the difference between the experimental and the calculated values is of interest. The difference per fluoride ion increases fairly systematically from left to right across the periodic table ruling out the possibility of there being a significant error in either the heat of formation of $\mathrm{F}^{-}$or the theoretical modeling of $\mathrm{F}^{-}$. In the case of $\mathrm{BF} \rightarrow \mathrm{B}^{+}+\mathrm{F}^{-}$, the calculation was repeated at the level of configuration interaction with all double substitutions. The energy difference changed by only $3 \mathrm{~kJ} \mathrm{~mol}^{-1}$, suggesting that the neglect of correlation energy is not important. On the other hand, the fact that the sign of the discrepancy is such that the molecule is calculated to be less stable with respect to the ions than is observed suggests that a more flexible basis set might result in better results at the SCF level. One obvious approach would to be to use scale factors optimized for each molecule, but as the results are considered sufficiently good for the purpose of predicting fluoride affinities, it did not appear that
the additional expense that this would entail was justified.
Table II lists the calculated geometries and energies of a number of species $\mathrm{X}=\mathrm{MF}_{n-p}^{p+}$; and in the last column, the fluoride affinity $\left(E_{\mathrm{F}}\right)$ of X , i.e., the energy change for the process $\mathrm{XF}^{-} \rightarrow \mathrm{X}+\mathrm{F}^{-}$. In determining this quantity, the calculated energy difference has been multiplied by the ratio experimental:theoretical from Table I on the assumption that the errors in the calculated values for a given sequence are systematic. This correction is very small anyway for the fluoride affinities of most current interest (such as that of $\mathrm{BF}_{3}$, for which it is only $3 \mathrm{~kJ} \mathrm{~mol}{ }^{-1}$ ).

The very good relative agreement between experimental and theoretical energies shown in Table I suggests that the calculated fluoride affinities should be sufficiently reliable to provide useful checks of reported values and reliable estimates of previously unknown ones.

## Comparison with Experiment

A number of fluoride affinities have been estimated from experimental data. They are usually reported as $\Delta H$ for the process

$$
\mathrm{XF}^{-} \rightarrow \mathrm{X}+\mathrm{F}^{-}
$$

(here X may be a neutral or a charged species). To compare these results with the $E_{\mathrm{F}}$ of Table II, there must be added to $E_{\mathrm{F}}$ (i) a correction for the change in zero point energy (typically $\sim 9 \mathrm{~kJ}$ $\mathrm{mol}^{-1}$ ) and (ii) a correction for heating from 0 K to the temperature at which $\Delta H$ is measured (usually $\sim 298 \mathrm{~K}$, but not always clear); this is $\sim-5 R T / 2 \sim-6 \mathrm{~kJ} \mathrm{~mol}^{-1}$. As the two corrections are in the opposite sense and their sum is less than typical uncertainties, in the following the calculated $E_{\mathrm{F}}$ is compared directly with reported values of $\Delta H^{\circ}$. Sometimes it is not clear what value of the electron affinity of fluorine or other reference data are assumed; this introduces further difficulties into comparison of theory and experiment.
(a) $\mathrm{BF}_{3}$. The fluoride affinity of $\mathrm{BF}_{3}$ is important to the experimental evaluation of other affinities. Reported values range from $259 \pm 40 \mathrm{~kJ} \mathrm{~mol}^{-1}(\mathrm{BH})^{4}$ to $385 \pm 25 \mathrm{~kJ} \mathrm{~mol}^{-1}(\mathrm{BH}),{ }^{5}$ with $297 \mathrm{~kJ} \mathrm{~mol}^{-1}$ (ICR) also being quoted. ${ }^{3}$ Mallouk et al. ${ }^{5}$ make a strong case for the higher value being preferred. The value of $E_{\mathrm{F}}$ reported here ( $358 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ) supports that contention.
(b) $\mathbf{S F}_{4}$. Murphy and Beauchamp ${ }^{2}$ estimate for the fluoride affinity of $\mathrm{SF}_{4}$ the value of $226 \pm 50 \mathrm{~kJ} \mathrm{~mol}^{-1}$ (ICR) and Larson and $\mathrm{McMahon}^{1 \mathrm{a}}$ give $183 \mathrm{~kJ} \mathrm{~mol}^{-1}$ (ICR). The higher value is in better agreement with $E_{\mathrm{F}}=211 \mathrm{~kJ} \mathrm{~mol}^{-1}$ reported here.
(c) $\mathrm{PF}_{5}$. Larson and McMahon ${ }^{1}$ give for the fluoride ion affinity of $\mathrm{PF}_{5}$ the value $356 \pm 40 \mathrm{~kJ} \mathrm{~mol}^{-1}$ (ICR). Mallouk et al., ${ }^{5}$ on the other hand, prefer $423 \pm 33 \mathrm{~kJ} \mathrm{~mol}^{-1}(\mathrm{BH})$. The latter value is closer to the calculated $E_{\mathrm{F}}=406 \mathrm{~kJ} \mathrm{~mol}^{-1}$.
(d) $\mathbf{S F}_{3}{ }^{+}$. Mallouk et al. ${ }^{5}$ quote $883 \pm 33 \mathrm{~kJ} \mathrm{~mol}^{-1}$ for the fluoride affinity of $\mathrm{SF}_{3}{ }^{+}(\mathrm{BH})$. This is very close to the calculated $E_{\mathrm{F}}=885 \mathrm{~kJ} \mathrm{~mol}^{-1}$.
(e) $\mathrm{PF}_{3}$. Larsen and McMahon ${ }^{1}$ give $168 \mathrm{~kJ} \mathrm{~mol}^{-1}$ (ICR) compared with a calculated $E_{\mathrm{F}}=189 \mathrm{~kJ} \mathrm{~mol}^{-1}$ for $\mathrm{PF}_{3}$.
(f) $\mathrm{SiF}_{4}$. The same authors give $251 \pm 20 \mathrm{~kJ} \mathrm{~mol}^{-1}$ (ICR) compared with a calculated $E_{\mathrm{F}}=323 \mathrm{~kJ} \mathrm{~mol}^{-1}$ for $\mathrm{SiF}_{4}$. On the
other hand, Jenkins and Pratt ${ }^{16}$ calculate (BH) for the double electron affinity of $\mathrm{SiF}_{4}\left(\rightarrow \mathrm{SiF}_{6}{ }^{2-}\right) \Delta E=138 \mathrm{~kJ} \mathrm{~mol}^{-1}$ in excellent agreement with the value ( $136 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ) calculated from Table II.
(g) $\mathrm{CF}_{3}{ }^{+}$. Data from ref 13 give for $\Delta H^{\circ}{ }_{298}=1158 \mathrm{~kJ} \mathrm{~mol}^{-1}$ (provenance unknown) for $\mathrm{CF}_{4} \rightarrow \mathrm{CF}_{3}{ }^{+}+\mathrm{F}^{-}$. The calculated value is $E_{\mathrm{F}}=1136 \mathrm{~kJ} \mathrm{~mol}^{-1}$.
(h) $\mathrm{BF}_{2}{ }^{+}$. Data from ref 15 (based on appearance potentials) yield $\Delta H^{\circ}{ }_{298}=1227 \pm 12 \mathrm{~kJ} \mathrm{~mol}^{-1}$ for $\mathrm{BF}_{3} \rightarrow \mathrm{BF}_{2}{ }^{+}+\mathrm{F}^{-}$. The calculated $E_{\mathrm{F}}=1262 \mathrm{~kJ} \mathrm{~mol}^{-1}$.

The comparison above shows generally good agreement between calculated values and experimental fluoride affinities determined by using a Born-Haber cycle (some of the residual differences may well be traced to the use of different reference data) providing a useful confirmation of the validity of that technique. The agreement reported in (g) and (h) above is also very satisfactory. On the other hand the fluoride affinities reported in the ICR ${ }^{1-3}$ literature are generally lower than those calculated suggesting a possible systematic error due to an incorrect reference value.

Comparison of calculated and observed geometries show that for the neutral molecules distances are generally in agreement within $\pm 0.02 \AA$ and angles within $\pm 1.0^{\circ}$, so that the data provided here should also provide useful estimates of the dimensions of species not as yet studied experimentally.

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Registry No. $\mathrm{F}^{-}$, 16984-48-8; HF, 7664-39-3; $\mathrm{HF}_{2}{ }^{-}$, 18130-74-0; $\mathrm{Li}^{+}$, 17341-24-1; $\mathrm{LiF}, 7789-24-4 ; \mathrm{LiF}_{2}^{-}, 50561-24-5 ; \mathrm{Be}^{2+}, 22537-20-8 ; \mathrm{BeF}^{+}$, 33808-13-8; $\mathrm{BeF}_{2}, 7787-49-7 ; \mathrm{BeF}_{3}{ }^{-}, 19181-26-1 ; \mathrm{BeF}_{4}{ }^{2-}$, 18539-20-3; Be , 7440-41-7; $\mathrm{BeF}^{-}, 51233-20-6 ; \mathrm{B}^{3+}, 22537-21-9 ; \mathrm{BF}^{2+}, 86686-93-3 ; \mathrm{BF}_{2}{ }^{+}$, 12355-90-7; $\mathrm{BF}_{3}, 7637-07-2 ; \mathrm{BF}_{4}^{-}, 14874-70-5 ; \mathrm{B}^{+}, 14594-80-0 ; \mathrm{BF}$, 13768-60-0; $\mathrm{BF}_{2}^{-}, 26202-31-3 ; \mathrm{C}^{4+}, 16092-62-9 ; \mathrm{CF}^{3+}, 102519-41-5$; $\mathrm{CF}_{2}{ }^{2+}$, 97411-56-8; $\mathrm{CF}_{3}{ }^{+}, 18851-76-8 ; \mathrm{CF}_{4}, 75-73-0 ; \mathrm{C}^{2+}, 16092-61-8$; $\mathrm{CF}^{+}$, 2154-59-8; $\mathrm{CF}_{2}, 2154-59-8 ; \mathrm{CF}_{3}^{-}, 54128-17-5 ; \mathrm{N}^{3+}, 18396-22-0$; $\mathrm{NF}^{2+}, 102519-42-6 ; \mathrm{NF}_{2}{ }^{+}, 31685-31-1 ; \mathrm{NF}_{3}, 7783-54-2 ; \mathrm{Na}^{+}, 17341-$ 25-2; $\mathrm{NaF}, 7681-49-4$; $\mathrm{NaF}_{2}{ }^{-}$, 50417-28-2; $\mathrm{Mg}^{2+}, 22537-22-0 ; \mathrm{MgF}^{+}$, 21308-25-8; $\mathrm{MgF}_{2}$, 7783-40-6; $\mathrm{MgF}_{3}{ }^{-}, 55161-80-3 ; \mathrm{MgF}_{4}{ }^{2-}$, 54097-16-4; $\mathrm{Mg}, 7439-95-4 ; \mathrm{MgF}^{-}, 39343-43-6 ; \mathrm{Al}^{3+}, 22537-23-1 ; \mathrm{AlF}^{2+}, 21330-18-7$; $\mathrm{AlF}_{2}{ }^{+}, 21559-03-5 ; \mathrm{AlF}_{3}, 7784-18-1 ; \mathrm{AlF}_{4}^{-}, 21340-02-3 ; \mathrm{AIF}_{5}^{2-}, 28747-$ 75-3; $\mathrm{AlF}_{6}{ }^{3-}, 21340-03-4 ; \mathrm{Al}^{+}, 14903-36-7$; AlF, 13595-82-9; $\mathrm{AlF}_{2}{ }^{-}$, 12758-11-1; $\mathrm{Si}^{4+}, 22537-24-2 ; \mathrm{SiF}^{3+}, 102519-43-7 ; \mathrm{SiF}_{2}{ }^{2+}, 102519-44-8$; $\mathrm{SiF}_{3}{ }^{+}, 38192-99-3 ; \mathrm{SiF}_{4}, 7783-61-1 ; \mathrm{SiF}_{5}^{-}, 21476-45-9 ; \mathrm{SiF}_{6}{ }^{2-}, 17084-$ 08-1; $\mathrm{Si}^{2+}, 14175-55-4 ; \mathrm{SiF}^{+}, 12518-09-1 ; \mathrm{SiF}_{2}, 13966-66-0 ; \mathrm{SiF}_{3}{ }^{-}$, 54711-93-2; $\mathrm{P}^{5+}, 22537-25-3 ; \mathrm{PF}_{2}{ }^{3+}, 102519-45-9 ; \mathrm{PF}_{3}{ }^{2+}, 102519-46-0$; $\mathrm{PF}_{4}{ }^{+}, 29075-80-7 ; \mathrm{PF}_{5}, 7647-19-0 ; \mathrm{PF}_{6}{ }^{-}, 16919-18-9 ; \mathrm{P}^{3+}, 20337-88-6$; $\mathrm{PF}^{2+}, 11080-57-2 ; \mathrm{PF}_{2}{ }^{+}, 37366-67-9 ; \mathrm{PF}_{3}, 7783-55-3 ; \mathrm{PF}_{4}^{-}, 25443-47-4 ;$ $\mathrm{S}^{6+}, 22537-26-4 ; \mathrm{SF}_{3}{ }^{3+}, 102519-47-1 ; \mathrm{SF}_{4}{ }^{2+}, 69754-88-7 ; \mathrm{SF}_{5}{ }^{+}$, 19167-14-7; $\mathrm{SF}_{6}, 2551-62-4 ; \mathrm{S}^{4+}, 20681-10-1 ; \mathrm{SF}_{3}{ }^{+}, 25431-36-1 ; \mathrm{SF}_{4}, 7783-$ 60-0; $\mathrm{SF}_{5}^{-}$, 31140-82-6.
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